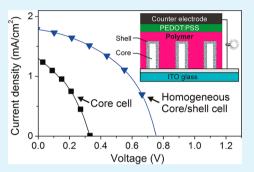
# Improved Open-Circuit Voltage in Polymer/Oxide-Nanoarray Hybrid Solar Cells by Formation of Homogeneous Metal Oxide Core/Shell Structures

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**Supporting Information** 

**ABSTRACT:** Incorporation of vertically aligned nanorod/nanowire arrays of metal oxide (oxide-NAs) with a polymer can produce efficient hybrid solar cells with an ideal bulk-heterojunction architecture. However, polymer/oxide-NAs solar cells still suffer from a rather low (normally, < 0.4 V) open-circuit voltage ( $V_{oc}$ ). Here we demonstrate, for the first time, a novel strategy to improve the  $V_{oc}$  in polymer/oxide-NAs solar cells by formation of homogeneous core/shell structures and reveal the intrinsic principles involved therein. A feasible hydrothermal—solvothermal combined method is developed for preparing homogeneous core/shell nanoarrays of metal oxides with a single-crystalline nanorod as core and the aggregation layer of corresponding metal oxide quantum dots (QDs) as shell, and the shell thickness (L) is easily controlled by the solvothermal reaction time for growing QDs on the nanorod. The core/shell



formation dramatically improves the device  $V_{oc}$  up to ca. 0.7–0.8 V depending on *L*. Based on steady-state and dynamic measurements, as well as modeling by space-charge-limited current method, it is found that the improved  $V_{oc}$  originates from the up-shifted conduction band edge in the core by the interfacial dipole field resulting from the decreased mobility difference between photogenerated electrons and holes after the shell growth, which increases the energy difference between the quasi-Fermi levels of photogenerated electrons in the core and holes in the polymer for a higher  $V_{oc}$ . Our results indicate that increasing  $V_{oc}$  by the core/shell strategy seems not to be dependent on the kinds of metal oxides.

**KEYWORDS:** solar cells, interface, nanoarray, metal oxide, electron mobility

## 1. INTRODUCTION

Hybrid polymer-based solar cells (HPSCs) that use conjugated polymers as electron donor (D) and inorganic semiconductor nanocrystals as electron acceptor (A) in a bulk-heterojunction (BHJ) architecture have shown promise for future photovoltaic devices.<sup>1-3</sup> Using nanoarrays (NAs) of vertically aligned nanorod/nanowire to replace the disordered A-phase pathways formed by nanoparticles can provide an ideal BHJ architecture for efficient solar cells. The metal oxide nanoarrays (oxide-NAs) (e.g., ZnO-NAs and TiO<sub>2</sub>-NAs) are often used to fabricate such aligned BHJ devices, because of their nontoxicity, high electron mobility and thermal stability.<sup>4-6</sup> However, the power conversion efficiency ( $\eta$ ) in polymer/ZnO-NA<sup>1,7–9</sup> or polymer/TiO<sub>2</sub>-NA<sup>10,11</sup> solar cells is not high yet (normally,  $\eta$  < 0.6%), even though  $\eta = 3.9\%$  has been achieved by incorporating PCBM into the HPSCs based on ZnO-NAs and poly(3-hexylthiophene) (P3HT);<sup>12</sup> in particular, the opencircuit voltage  $(V_{oc})$  in the polymer/oxide-NA devices is generally rather low (0.1-0.5 V), which significantly limits the device  $\eta$ . To establish the methods for achieving high  $V_{0c}$  and understand the principles therein are the challenges in access to efficient polymer/oxide-NA solar cells.

Modification at the D/A interface with organic molecules generally enhances the short-circuit current ( $J_{\rm sc}$ ) in polymer/

oxide-NAs devices by increasing the interfacial compatibility, but hardly improves  $V_{oc}$ , <sup>8,10,13-15</sup> Coating inorganic layers (e.g., TiO<sub>2</sub>)<sup>16</sup> TiO<sub>x</sub>, <sup>17</sup> CdS, <sup>18</sup> CdSe, <sup>19,20</sup> PbS<sup>21</sup>) on oxide-NAs to form heterogeneous structures has been demonstrated to increase significantly the  $V_{oc}$  in polymer/Oxide-NA solar cells. Deposition of crystalline TiO<sub>2</sub> or amorphous titanium oxide  $(TiO_x)$  on ZnO-NA provided ZnO/TiO<sub>2(x)</sub>-core/shell nanorod arrays (ZnO-TiO<sub>2</sub>-NAs<sup>16</sup> or ZnO-TiO<sub>x</sub>-NAs<sup>17</sup>), where the P3HT/ZnO-TiO<sub>2</sub>-NA solar cells had a  $V_{oc}$  up to 0.55 V and an increased  $J_{sc}$  depending on the TiO<sub>2</sub> shell thickness, whereas the P3HT/ZnO-TiO<sub>x</sub>-NA devices exhibited a much higher  $V_{oc}$ (~0.80 V) than P3HT/ZnO-NA devices but a much lower  $J_{sc}$ due to the barrier effect of  $TiO_x$  layer on the charge transfer from the polymer to ZnO;<sup>17</sup> sensitization of ZnO-NAs with CdS (or PbS) quantum dots (QDs) produced the  $V_{oc}$  up to 0.85 V (or 0.59 V) in their HPSCs, along with an increased  $J_{sc}$  to a certain extent.<sup>18,21</sup> Heterogeneous modification inevitably generates the interfacial defects due to lattice mismatch<sup>22</sup> or lattice expansion<sup>23</sup> to affect the charge transport and recombination. Moreover, the effects of the heterogeneous

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shell layers on device performance are hitherto not clearly understood yet. For examples, the improved  $V_{oc}$  by the TiO<sub>x</sub> shell formation is explained from the decrease in reverse saturation current based on a modified ideal diode equation;<sup>17</sup> the CdS-sensitization effect on  $V_{oc}$  is attributed to the band offset between CdS and ZnO in the cascading energy band structure formed upon CdS deposition.<sup>18</sup> However, it is hard to correlate the loads of the modifiers (e.g., CdS-QDs) with the device performance in those explanations.

The factors determining  $V_{oc}$  in HPSCs are still the subject of much debate. Normally, the energy difference between the highest occupied molecular orbital (HOMO) level of the donor and the conduction band edge of the acceptor determines the maximum  $V_{oc}^{1,24,25}$  Recent studies demonstrated that the change of  $V_{\rm oc}$  in HPSCs may correlate with the altered band alignment at the D/A interface due to the presence of interfacial dipoles.<sup>9,26–28</sup> For example, the improved  $V_{oc}$  in the P3HT/ZnO-NA device subjected to the UV/ozone treatment of ZnO surface is believed to originate from the band alignment changed by interfacial dipole formation between P3HT and  $ZnO_{3}^{9}$  the molecular dipoles have been used as modifier to shift the band positions of  $TiO_{2}^{26}$  or  $ZnO^{27,28}$  for adjusting  $V_{oc}$  in their HPSC devices. In addition, Blom and co-workers<sup>29,30</sup> found that the balance between the carrier extraction and the energy loss via charge recombination led to a distinct optimization in the carrier mobility for efficient fullerenebased devices, and too high mobility for charge extraction will lead to a low device  $V_{oc}$  but the process governing  $V_{oc}$  was not fully understood even though a correlation between a low  $V_{oc}$ and a reduced difference between the quasi-Fermi levels of electrons and holes was observed. There is still insufficient knowledge concerning the intrinsic correlation between  $V_{oc}$  and charge carrier mobility in HPSCs, whereas a few studies had been reported on the relationship between carrier mobility and  $J_{\rm sc}^{29,31}$ 

In this paper, we develop a novel and feasible strategy to improve the  $V_{oc}$  in polymer/oxide-NAs solar cells by formation of homogeneous core/shell nanoarrays. As depicted in Figure 1,

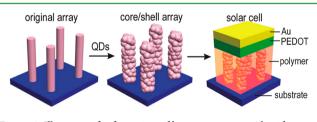


Figure 1. Illustration for formation of homogeneous metal oxide core/ shell nanoarrays and their solar cell architecture.

the hydrothermally preformed oxide-NAs are modified by their corresponding QDs through a solvothermal process, generating core/shell nanoarrays (referred to as oxide-NR-QD-NAs); then the core/shell nanoarrays are infiltrated with poly(2-methoxy-5-(2-ethylhexyloxy)-1,4-phenylene vinylene) (MEH-PPV) to fabricate solar cells. Results show that the high  $V_{\rm oc}$  subjected to the core/shell formation originates from the generation of interfacial dipoles due to the decreased mobility difference between photogenerated electrons and holes.

## 2. EXPERIMENTAL SECTION

Synthesis of ZnO-NR-QD-NAs. ZnO-NA was hydrothermally grown by suspending the ZnO-coated indium tin oxide (ITO) conducting glass (<15  $\Omega/\Box$ , Shenzhen Laibao Hi-Tech Co., Ltd.,

China) upside down in an aqueous solution of zinc nitrate hexahydrate (0.025 M) (Sinopharm Chemical Reagent Co., Ltd.) and hexamethylenetetramine (0.025 M) (99.99%, Alfa Aesar) at 90 °C as reported previously.<sup>32</sup> The ZnO-NA was put upside down in a Teflon-lined stainless steel autoclave (100 mL volume), and added 60 mL zinc acetate solution in ethanol (0.02 M); the sealed autoclave was heated at 95 °C for a defined time (t = 4, 4.5, or 6 h) for the solvothermal growth of ZnO-QDs, to get ZnO-NR-QD-NA with a specific shell thickness; finally, after the autoclave naturally cooled to room temperature, the ZnO-NR-QD-NA was taken out, thoroughly rinsed with alcohol and dried in a N<sub>2</sub> flow.

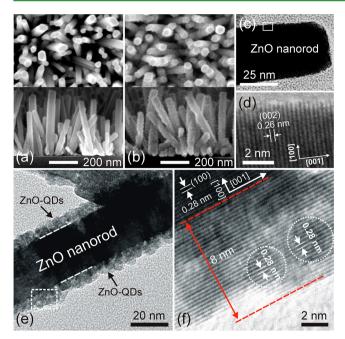
**Synthesis of TiO<sub>2</sub>-NR-QD-NAs.** Fluoride-doped tin oxide (FTO) coated glass sheet (14  $\Omega/\Box$ , 400 nm FTO in thickness, Nippon Sheet Glass Co.) was first patterned into stripes (12 × 4 mm<sup>2</sup>) on the glass substrate by HCl solution and Zn powder, and washed twice with acetone, isopropanol, and deionized water, respectively, in which the ultrasonic treatment for 10 min was done for each wash. TiO<sub>2</sub>-NA was hydrothermally grown on FTO surface.<sup>33</sup> After a thermal annealing in air at 450 °C for 30 min, the TiO<sub>2</sub>-NA substrate was put upside down in a Teflon-lined stainless steel autoclave (100 mL volume), and added 60 mL titanium isopropoxide solution in ethanol (0.1 M); the sealed autoclave was heated at 200 °C for 4 h for the solvothermal growth of TiO<sub>2</sub>-QDs, to get TiO<sub>2</sub>-NR-QD-NA. After the autoclave naturally cooled to room temperature, the substrate was taken out, thoroughly rinsed with alcohol and dried in a N<sub>2</sub> flow.

**Device Fabrication.** MEH-PPV (average  $M_p = 40000-70000$ , Aldrich) and poly(3,4-ethylene dioxythiophene):poly(styrenesulfonate) (PEDOT:PSS) (Clevios P HC V4, H. C. Starck) was commercially obtained. Solar cells were fabricated by the procedure similar to our previous reports,  $^{13,32,34}$  where the polymer layer was deposited on the top of nanoarray by spin-coating (1500 rpm, 40 s) the MEH-PPV solution in chlorobenzene (5 mg/mL) under ambient conditions for twice, followed by annealing at 150 °C under N<sub>2</sub> atmosphere for 10 min to ensure the polymer infiltration into the interspaces between nanorods. After the spin-coating of PEDOT:PSS (2000 rpm, 60 s), the sample was heated for 15 min at 100 °C under a nitrogen atmosphere. Finally, a gold electrode (100 nm) was evaporated through a shadow mask to form an overlapped area between ITO and Au of 3 mm  $\times$  3 mm, which defined as the effective device area. The electron-only devices were fabricated the procedure same to the solar cell assembly except for the PEDOT/Au top electrode was replaced with the evaporated LiF (1 nm)/Al (100 nm).

Characterization. X-ray diffraction (XRD) patterns were recorded on an MXP18AHF X-ray diffractometer with Cu K<sub> $\alpha$ </sub> irradiation ( $\lambda$  = 1.54056 Å). Scanning electron microscopy (SEM) measurements were carried out on a field-emission scanning electron microscopy (FE-SEM, FEI Sirion200). Transmission electron microscopy (TEM) and high-resolution TEM (HRTEM) studies were performed on a JEOL-2010 microscope under an acceleration voltage of 200 kV. Photoluminescence (PL) measurements were performed on a Hitachi F-7000 spectrofluorophotometer. Solar cells were characterized with current-voltage (I-V) measurements and intensity modulated photovoltage spectroscopy (IMVS). Steady-state J-V curves were measured under AM1.5 illumination (100 mW/cm<sup>2</sup>, calibrated with a standard crystalline silicon solar cell) under ambient conditions on a 94023A Oriel Sol3A solar simulator (Newport Stratford, Inc.) with a 450 W xenon lamp as the light source, and the J-V data were collected with an Oriel I-V test station (PVIV-1A, Keithley 2400 Source Meter, Labview 2009 SP1 GUI Software, Newport). IMVS measurements were done on a controlled intensity modulated photo spectroscopy (CIMPS) (Zahner Co., Germany) in ambient conditions, as described previously.<sup>32,34</sup> The J-V curves of the solar cells and the electron-only devices in the dark were collected on the CIMPS system with ITO contact as negative electrode.

## 3. RESULTS AND DISCUSSION

**3.1. Formation of Core/Shell ZnO Nanoarrays.** SEM image (Figure 2a) shows that the initial ZnO-NA is vertically aligned on ITO substrate, with a length of ~400 nm, a diameter

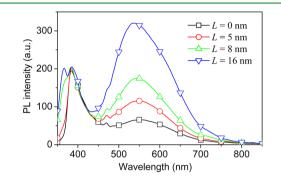


**Figure 2.** SEM images of ZnO-NA (a) before and (b) after deposition of ZnO-QDs. (c, e) TEM and (d, f) HRTEM images of (c, d) ZnO nanorod and (e, f) ZnO-NR-QD nanorod with L = 8 nm, in which the ZnO and ZnO-NR-QD nanorods were obtained from the samples in a and b, respectively. The HRTEM images were taken from the marked regions on the corresponding TEM images, and the cycles in (f) identify individual ZnO-QDs.

of 20–50 nm and a number density of about  $4-5 \times 10^2/\mu m^2$ , agreeing with previous results;<sup>32,35</sup> after deposition of ZnO-QDs, the morphology of ZnO-NA is not remarkably altered, but ZnO nanorod surfaces become decorated by small particles (Figure 2b). The ZnO-NA was indexed to wurtzite ZnO and the single-crystalline ZnO nanorods therein grew along [001] crystallographic direction,<sup>35,36</sup> as indicated by XRD pattern (see Figure S1 in the Supporting Information). TEM image (Figure 2c) shows that the surface of original ZnO nanorod is quite smooth, and HRTEM image (Figure 2d) confirms the ZnO nanorod growth along [001] direction. After the ZnO-QDs growth, the ZnO nanorod surface is coated with a shell layer (Figure 2e), and the XRD pattern of the core/shell sample (see Figure S1 in the Supporting Information) was only indexed to wurtzite ZnO. As shown in Figure 2f, both core and shell exhibit the same oriented lattice fringes without lattice mismatch or distortion, where the marked interplanar spacings of 0.28 nm correspond to the (100) lattice plane of wurtzite ZnO, suggesting that the shell consisting of ZnO-QDs with a size of 2-4 nm was actually epitaxially grown on ZnO nanorod surface. The aggregation of epitaxially grown ZnO-QDs on ZnO nanorods is similar to the growth of ZnO nanoparticles on ZnS nanobelt.<sup>37</sup> Both TEM and XRD results clearly demonstrate that, with the preformed ZnO-NA as a template, the ZnO-NR-QD-NA core/shell nanoarray with a singlecrystalline ZnO nanorod serves as core and an aggregation layer of ZnO-QDs as shell has been successfully prepared by the combined hydrothermal-solvothermal technique. Several ZnO-NR-QD-NAs with different shell thickness (L) were prepared in this experiment by controlling the solvothermal reaction time (t) for growing ZnO-QDs. With the reaction times of t = 4, 4.5, and 6 h, the L = 5, 8, and 16 nm were obtained, respectively (see Figure S2 in the Supporting

Information). In comparison to the original ZnO nanorods (Figure 2d), more irregular terminations of lattice periodicity appear inside the shell and on the shell surface (Figure 2f), indicating a higher concentration of defects is introduced into the shell.<sup>38</sup>

With increasing L (i.e., reaction time t), more surface defects are produced in ZnO-NR-QD-NAs, as is characterized by photoluminescence (PL) spectra (Figure 3). The ZnO-NA

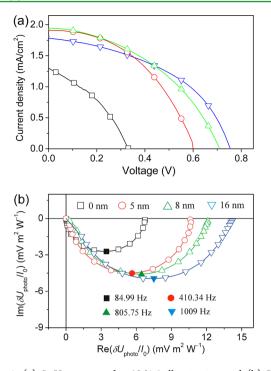


**Figure 3.** Room temperature PL spectra of original ZnO-NA (L = 0 nm) and ZnO-NR-QD-NAs with L = 5, 8, and 16 nm.

sample exhibits a sharp emission at 382 nm and a broad emission peak at 500-600 nm. The emission at 382 nm is regarded as near band edge (NBE) emission, and the broad emission is attributed to the deep level emission (DLE) due to intraband surface defects,  $^{39}$  such as oxygen vacancy (V\_0), zinc vacancy  $(V_{Zn})$ , oxygen interstitial  $(O_i)$ , zinc interstitial  $(Zn_i)$ , and oxygen antisites  $(O_{Zn})$ .<sup>40</sup> In the ZnO-NR-QD-NAs, the similar NBE emission and DLE peaks are observed, where the NBE emission intensity is almost unchanged but the DLE intensity increases dramatically as L increases, suggesting that the NBE emission is mainly from the ZnO nanord core and the increased DLE is related to the shell formation. The intensity ratio between the DLE and NBE emission is an indication of crystal quality, and a larger ratio means a higher concentration of surface defects.<sup>39</sup> Since the intensity of NBE emission in all the samples is quite similar, the higher DLE intensity in ZnO-NR-QD-NAs with increasing L actually indicates a higher defect concentration. In addition, an emission at 366 nm is observed and its intensity gets increased with increasing L, which has been recognized as the NBE emission of ZnO-QDs with a size less than 5 nm.<sup>41</sup> Hence, the appearance of the NBE peak at 366 nm further confirms the formation of shell consisting of ZnO-QDs in the ZnO-NR-QD-NAs.

**3.2. Solar Cells.** *3.2.1. Device Performance.* Figure S3 (Supporting Information) shows that MEH-PPV can compactly fill the interspaces between either ZnO-NA or ZnO-NR-QD-NA nanorods after thermal annealing (150 °C, 20 min), which is the same as the observation in our previous results<sup>13,32,34</sup> and the reports on the polymer infiltration in ZnO-TiO<sub>2</sub>-NAs<sup>16</sup> or CdS-QDs and CdSe-QDs cosensitized ZnO-NAs<sup>42</sup> by others. To investigate the device photovoltaic performance related to formation of core/shell structure, we compare *J*–*V* behaviors of the devices based on the as-synthesized ZnO-NA and ZnO-NR-QD-NAs under simulated AM1.5 illumination (100 mW/cm<sup>2</sup>) (Figure 4a). The averaged overall photovoltaic performance of three individual devices for each sample is presented in Table 1. The MEH-PPV/ZnO-NA device exhibits a rather low V<sub>oc</sub> (<0.4 V), similar to the previous reports on P3HT/ZnO-NA<sup>8,9,43</sup> and MEH-PPV/ZnO-NA.<sup>13</sup> In contrast, the ZnO-QDs shell





**Figure 4.** (a) J-V curves under AM1.5 illumination and (b) IMVS spectra of the solar cells based on ZnO-NA ( $\Box$ ) and ZnO-NR-QD-NAs with L = 5 nm ( $\bigcirc$ ), 8 nm ( $\triangle$ ), and 16 nm ( $\bigtriangledown$ ). The solid symbols in b identify the  $f_{\min}$  points.

formation improves significantly  $J_{sc}$  and  $V_{oc}$  depending on L. The device  $V_{oc}$  increases rapidly from 0.36 V for L = 0 nm to 0.74 V for L = 8 nm; further increasing L up to ~16 nm only leads to a very slight increase in  $V_{oc}$ . However, the device  $J_{sc}$ first increases when L increases from 0 to 8 nm, and then decreases for further increasing L, which is similar to the previously observed  $J_{sc}$  dependences on the TiO<sub>2</sub> shell thickness in P3HT/ZnO-TiO<sub>2</sub>-NA solar cells<sup>15</sup> and the amount of CdS-QDs in CdS-sensitized MEH-PPV/ZnO-NA devices.<sup>18</sup>

Here we first discuss the *L*-dependence of  $J_{\rm sc}$ . Our previous theoretical simulations revealed that the photocurrent in MEH-PPV/ZnO-NAs solar cells originates dominantly from the exciton dissociation at the MEH-PPV/ZnO interfaces formed by the side faces of ZnO nanorods.<sup>32</sup> Because of the short diffusion length of excitons in MEH-PPV (~5–8 nm),<sup>44–46</sup> only the excitons within the exciton diffusion length of the MEH-PPV/ZnO interface lead to free charge carriers. The interspacings between the ZnO nanorods in the as-synthesized ZnO-NA are not uniform and distributed from tens to hundred nanometers (Figure 2a), and mostly in the range of 5–50 nm consistent with Yang's report.<sup>35</sup> Statistically, the interspacing between the ZnO nanorods in the ZnO-NA is averaged to be  $25~\text{nm.}^{35,47}$  On the basis of the averaged ZnO nanorod interspacing in ZnO-NA and the diffusion length of excitons in MEH-PPV, it is theoretically predicted<sup>47</sup> that an optimized shell thickness  $(L_{opt})$  (regardless of the shell composition) of 4–8 nm exists for the efficient charge generation, and the polymer contribution to  $J_{sc}$  will not be influenced for the shell with a thickness of  $L \leq L_{opt}$  but be reduced as  $L > L_{opt}$ . The presence of the peak  $J_{sc}$  at L = 8 nm agrees well with our theoretical prediction. Moreover, formation of the coarse shell will inevitably increase the interfacial area for exciton dissociation in comparison to the smooth surface of the original ZnO nanorods. The increased  $J_{sc}$  with increasing L for  $L \leq 8$  nm is explained by the increased interfacial area for exciton dissociation with the shell formation, when the polymer contribution is not affected. Reasonably, with further increasing L up to the value with  $L > L_{opt}$ , the device  $J_{sc}$  will be a competitive result of increased interfacial area for exciton dissociation and decreased polymer amount infiltrated into nanorod interspaces. The decreased  $J_{sc}$  upon increasing L up to 16 nm is attributed to the reduced the amount of MEH-PPV between the nanorods,<sup>16,18,47</sup> which leads to a reduced MEH-PPV contribution to  $J_{sc}$  that suppresses the contribution from the increased interfacial area for charge generation.

On the other hand, shell formation increases the fill factor (FF) of the solar cells (Table 1), from 36% for L = 0 nm to 45% for L = 5 nm for example, the increase in FF due to the shell formation is attributed to the balanced mobility between photogenerated electrons and holes<sup>4,48,49</sup> as a result of the reduced mobility difference between the charge carriers (refer to later discussion). Furthermore, the shell thickness has a faint influence on the FF (around 45%) in the MEH-PPV/ZnO-NR-QD-NAs solar cells at the first glance (Table 1), but the detailed inspection shows that the FF indeed reduces with increasing *L*, even though slightly. The similar *L*-dependence of FF value was also observed in the P3HT/ZnO-TiO<sub>2</sub>-NAs devices,<sup>16</sup> for which the reason was not unclear. The slight decrease in FF upon increasing *L* in this experiment may be due to the reduced shunt resistance ( $R_{\rm sh}$ ) in device (Table 1).<sup>50</sup>

IMVS is a powerful dynamic photoelectrochemical method to characterize the electron lifetime  $\tau_{\rm e}$  related to interfacial charge recombination.<sup>13,34</sup> IMVS measures the photovoltage  $(\delta U_{\rm photo})$  response to a small sinusoidal light perturbation superimposed on the background light intensity  $I_0$  under opencircuit condition. The measured IMVS responses are shown in Figure 4b, and the observation that a larger  $V_{\rm oc}$  is accompanied by a higher photovoltage agrees with previous results.<sup>13,34</sup> The  $\tau_{\rm e}$  value was calculated from the frequency  $(f_{\rm min})$  of the lowest imaginary component of IMVS response according to the relation  $\tau_{\rm e} = (2\pi f_{\rm min})^{-1}$  and collected in Table 1. The  $\tau_{\rm e}$  values (0.15–1.75 ms) are comparable to our previous reports.<sup>13,34</sup>

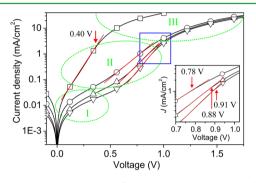
Table 1. Device Performance of MEH-PPV/ZnO-NA (L = 0 nm) and MEH-PPV/ZnO-NR-QD-NAs (L = 5, 8, and 16 nm) Hybrid Solar Cells under AM1.5 Illumination (100 mW/cm<sup>2</sup>)<sup>*a*</sup>

L (nm)	$V_{\rm oc}~({ m V})$	$J_{\rm sc}~({\rm mA/cm}^2)$	FF (%)	η (%)	$\tau_{\rm e}~({\rm ms})$	$R_{\rm s}~({ m k}\Omega)$	$R_{\rm sh}~({ m k}\Omega)$
0	$0.36 \pm 0.04$	$1.37 \pm 0.05$	$35.95 \pm 1.17$	$0.17 \pm 0.02$	$1.75 \pm 0.22$	$1.22 \pm 0.03$	$5.28 \pm 0.32$
5	$0.63 \pm 0.05$	$1.69 \pm 0.25$	$44.92 \pm 3.46$	$0.46 \pm 0.08$	$0.36 \pm 0.05$	$1.33 \pm 0.19$	$62.73 \pm 0.16$
8	$0.74 \pm 0.03$	$1.89 \pm 0.11$	$44.23 \pm 1.60$	$0.62 \pm 0.03$	$0.21 \pm 0.03$	$1.19 \pm 0.23$	$34.31 \pm 3.54$
16	$0.76 \pm 0.01$	$1.80 \pm 0.16$	$43.66 \pm 2.38$	$0.59 \pm 0.07$	$0.15 \pm 0.02$	$1.07 \pm 0.05$	$27.06 \pm 0.10$

"Each of the photovoltaic performance data with standard deviations represents the average of three devices.  $R_s$  and  $R_{sh}$  link with the slope characteristics at  $V_{oc}$  and  $J_{sc}$ , respectively.

recombination rate in the sample with a thicker shell. HRTEM and PL results have clearly indicated that a higher concentration of surface defects is introduced as *L* increases (Figures 2f and 3). The smaller  $\tau_e$  for the thicker shell is reasonably due to the increased defect concentration in shell, because as more electrons are trapped in surface defects they will not easily escape from the MEH-PPV/ZnO-QDs interface where a substantial number of holes exist, resulting in a faster recombination or smaller  $\tau_e$ .<sup>13,43</sup> Clearly, the *L*-dependence of  $\tau_e$  is opposite to that of  $V_{oc}$  (Table 1), which will be discussed in the later section.

As shown in Figure 5, the semilogarithmic plots of the J-V characteristics of solar cells in the dark normally exhibit three



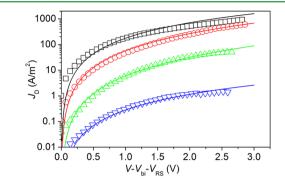
**Figure 5.** Semilogarithmic plots of J-V characteristic in the dark of the solar cells based on ZnO-NA ( $\Box$ ) and ZnO-NR-QD-NAs with L = 5 nm ( $\bigcirc$ ), 8 nm ( $\triangle$ ), and 16 nm ( $\bigtriangledown$ ). The inset shows the magnified J-V curve in the marked frame, to identify the  $V_{\rm bi}$  values in the core/shell samples.

distinct regimes,<sup>28,51</sup> that is, (I) linear increase for leakagedominated current, (II) exponential increase for spacecharge-limited current. The built-in voltage ( $V_{\rm bi}$ ) corresponding to the built-in electric filed ( $E_{\rm bi}$ ) due to the work-function ( $W_{\rm f}$ ) difference between ITO and Au counter electrodes<sup>34,52</sup> was evaluated at the point where the dark J-V curve begins to follow a quadratic behavior.<sup>28,51</sup> The measured  $V_{\rm bi}$  in MEH-PPV/ZnO-NA device is 0.40 V, but much higher  $V_{\rm bi}$  is produced with increasing L in MEH-PPV/ZnO-NR-QD-NA devices (Figure 5).

3.2.2. Formation of Dipoles. In MEH-PPV/ZnO-NA solar cells, the charge generation mainly originates from the exciton generated between the nanorods with a length of about 400 nm, in which the excitons mainly dissociate at nanorod side faces.<sup>32</sup> Hence, the behaviors of charge carriers nearby the ZnO nanorod side surface are very important for device performance. After the electron injecting into the ZnO-QDs shell, they will be either injected into the ZnO nanorod core or trapped by the abundant surface defects in the shell, where trapping process will inevitably reduce the electron mobility ( $\mu_e$ ). The  $\mu_e$ values in MEH-PPV/ZnO-NA and MEH-PPV/ZnO-NR-QD-NA composite films were obtained by space-charge-limited current (SCLC) method,  $^{53-55}$  using the electron-only devices fabricated by sandwiching the photoactive layers between ITO and LiF/Al electrodes (see Figure S4 in the Supporting Information). The SCLC with field-dependent mobility is given by<sup>55</sup>

$$J_{\rm e} = \frac{9}{8} \varepsilon_{\gamma} \varepsilon_0 \mu_{\rm e} \frac{V^2}{L^3} \exp\left(0.891 \gamma_{\rm e} \sqrt{\frac{V}{L_{\rm p}}}\right) \tag{1}$$

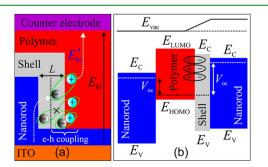
where  $J_{\rm e}$  is the electron current,  $\mu_{\rm e}$  the zero-field mobility of electron,  $\gamma_{\rm e}$  a field activation factor,  $\varepsilon_0$  the permittivity of free space,  $\varepsilon_{\rm r}$  the relative permittivity of the material, and  $L_{\rm p}$  the thickness of photoactive layer. The J-V characteristics in the dark  $(J_{\rm D}-V)$  of the electron-only devices were fitted to eq 1 to get parameter  $\mu_{\rm e}$  (Figure 6). Results show that the shell



**Figure 6.**  $J_{\rm D}-V$  of ITO/ZnO:MEH-PPV/LiF/Al electron-only devices with ZnO-NA ( $\Box$ ) and ZnO-NR-QD-NAs ( $\bigcirc$  5 nm,  $\triangle$  8 nm,  $\bigtriangledown$  16 nm). The solid lines imposed on scatters represent the fits of the experimental data (scatters) using eq 1. The applied voltage (V) corrected for the built-in voltage ( $V_{\rm bi}$ ) due to the work-function difference between Al and ITO electrodes and the voltage drop ( $V_{\rm RS}$ ) arising from the resistance of collection electrode. The fits produce the  $\mu_{\rm e}$  values of  $1.8 \times 10^{-6}$  ( $\Box$ ),  $3.9 \times 10^{-7}$  ( $\bigcirc$ ),  $3.7 \times 10^{-8}$  ( $\triangle$ ), and 2.8  $\times 10^{-9}$  ( $\bigtriangledown$ ) m<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> in the MEH-PPV/ZnO composite films.

formation leads to the  $\mu_e$  about 1–3 orders of magnitude lower than that in MEH-PPV/ZnO-NA film, with a smaller  $\mu_e$  for a larger *L*. For example,  $\mu_e$  values are 1.8 × 10<sup>-6</sup> and 3.7 × 10<sup>-8</sup> m<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> for MEH-PPV/ZnO-NA and MEH-PPV/ZnO-NR-QD-NA with *L* = 8 nm, respectively. The reduced  $\mu_e$  upon the shell formation is reasonably resulted from the reduced electron mobility in the shell layer, because the polymer and ZnO nanorod core remain unchanged in the samples.

Clearly, the electrons (e) in the shell and the holes (h) in the polymer in MEH-PPV/ZnO-NR-QD-NA devices have a smaller mobility difference (the hole mobility is  $1 \times 10^{-10}$  m<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> in MEH-PPV<sup>56,57</sup>) than in MEH-PPV/ZnO-NA ones, which will favor the electronic coupling between the electrons and holes to form dipoles directing toward the polymer (negative charge in ZnO-QDs shell and positive charge in polymer), as depicted in Figure 7a; in addition, the dipoles orient along the built-in filed  $E_{\rm bi}$  direction toward the counter electrode.<sup>28</sup> Note that, a certain number of dipoles can also be



**Figure 7.** (a) Illustration of interfacial dipole formation, the cross between electron and hole means they are not able to couple with each other. (b) The band structure changes subjected to the core/shell formation, where the dashed lines indicate the quasi-Fermi levels.

produced at the polymer/ZnO interface by the electrons trapped on ZnO nanorods and the hole in the polymer in the MEH-PPV/ZnO-NA system,<sup>58</sup> but formation of core/shell structure eventually increases the dipole number.

Obviously, the photogenerated electrons and holes need to locate within a certain e-h separation where they are spatially close enough to couple with each other (ca.  $3-10 \text{ nm}^{59}$ ) for the dipole formation. Therefore, the e-h coupling region for dipole formation is very thin (less than 10 nm) and the dipole formation inevitably correlates with the shell thickness L, resulting in the presence of an expected maximum L of around 10 nm, above which dipole formation will become saturated.<sup>60</sup> As the shell thickness is comparable to the e-h separation for coupling, all the electrons in the shell are capable to form dipoles; however, in the shells with a larger thickness (e.g., L >10 nm), some injected electrons will diffuse toward the core with a spatial separation beyond the e-h coupling region (Figure 7a) and they will not able to couple with the holes in the polymer to generate dipoles, eventually causing the presence of saturated electron concentration for coupling.

The interfacial dipole shell (directing toward the polymer) can induce an extra polar electric field  $(E'_{bi})$  as an addition to enhance  $E_{\rm bi}$  (Figure 7a),<sup>28,61</sup> this is confirmed by the observed *L*-dependence of  $V_{\rm bi}$ . According to the  $W_{\rm f}$  difference between Au (~5.1 eV)<sup>6</sup> and ITO (~4.8 eV),<sup>6</sup> the maximum theoretical  $V_{\rm bi}$  in the solar cells is estimated to be about 0.30 V. However, the measured V<sub>bi</sub> in MEH-PPV/ZnO-NA device is actually 0.40 V, and much higher  $V_{\rm bi}$  (0.78–0.91 V) is produced with increasing L in MEH-PPV/ZnO-NR-QD-NA devices (L = 5-16 nm) (Figure 5). All the measured  $V_{\rm bi}$  is significantly larger than the theoretical expectation (0.30 V), suggesting the extra voltage contribution to  $V_{\rm bi}$  from the dipoles present at the polymer/ZnO interface.<sup>60</sup> The larger  $V_{\rm bi}$  in MEH-PPV/ZnO-NR-QD-NA solar cells than in MEH-PPV/ZnO-NA devices demonstrates more dipoles generated due to the shell formation. The L-dependence of V<sub>bi</sub> indicates a L-dependent dipole formation in the ZnO-QDs shell; in particular, increasing L from 8 to 16 nm leads to a very slight increase in  $V_{\rm bi}$  from 0.88 to 0.91 V, indicating the dipole formation has almost become saturated at  $L \ge 8$  nm. The observed maximum L of around 8 nm for increasing  $V_{\rm bi}$  agrees with the expectation on the dipole formation.

3.2.3. Dipole Effects on  $V_{oc}$  and  $\tau_e$ . Quantum confinement effect will shift the conduction band  $(E_c)$  and valence  $(E_v)$ edges to increase in band gap of QDs with reducing size. Hence, there exists a band offset between core and shell in ZnO-NR-QD-NAs (Figure 7b), which may increase the energy difference between the HOMO level of the polymer and the  $E_c$ edge of the QDs acceptor to affect the  $V_{oc}$ .<sup>18,21,62</sup> In this experiment, the increased  $V_{\rm oc}$  upon increasing L is reasonably not resulted from the upshift of  $E_c$  edge in ZnO-QDs shell due to the quantum confinement effect. On one hand, the  $E_c$  shift of ZnO-QDs relative to bulk value is less than 0.13 eV when the particle size is larger than 3 nm (see Figure S5 in the Supporting Information). However, the observed improvement of eV<sub>oc</sub> in ZnO-NR-QD-NA devices are much larger than the expected maximum from the  $E_c$  shift; moreover, a thicker shell inevitably contains larger ZnO-QDs (Figure 2f and Figure S2 in the Supporting Information), the ZnO-NR-QD-NA device with a higher  $V_{\rm oc}$  should be observed for thinner shell based on the band-offset standpoint due to a larger energy difference between the polymer HOMO level and the  $E_c$  of smaller ZnO-QDs. Therefore, the upshift of the  $E_c$  edge in shell due to

the quantum confinement effect is not the reason for the increased  $V_{\rm oc}.$ 

The L-dependences of  $V_{\rm oc}$ ,  $V_{\rm bi}$ ,  $\tau_{\rm e}$ , and  $\mu_{\rm e}$  are depicted in Figure 8. The  $\mu_{\rm e}$  values keep a remarkable decrease with

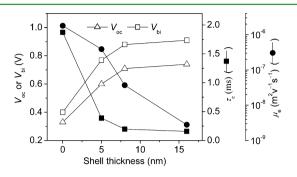


Figure 8. Dependences of  $V_{oc}$ ,  $V_{bi}$ ,  $\mu_{e}$ , and  $\tau_{e}$  on ZnO shell thickness.

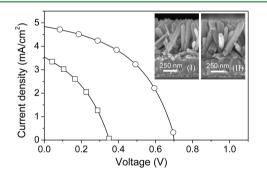
increasing *L* from 0 to 16 nm, for which the increasing surface defect concentration in the growing shell is responsible; however,  $V_{oc}$  increases initially when increasing *L* from 0 to 8 nm and almost remains unchanged with further increasing *L* up to 16 nm. The different *L*-dependences of  $V_{oc}$  and  $\mu_e$  infer there is no direct relation between  $V_{oc}$  and  $\mu_{e}$  in other words,  $V_{oc}$  is not directly determined by  $\mu_e$ . In contrast,  $V_{bi}$  and  $V_{oc}$  follow the same changing trend with increasing *L*, indicating a strong correlation between  $V_{oc}$  and dipole formation in the shell.<sup>51,60</sup>

It is believed that the Voc in the MEH-PPV/ZnO-NA and MEH-PPV/ZnO-NR-QD-NA devices is mainly determined by the energy levels of  $E_c$  edge in ZnO nanorod core and HOMO band in MEH-PPV, but significantly correlates with the quasi-Fermi levels of the electrons in ZnO nanorod core and the holes in polymer.<sup>13,61,63</sup> With the presence of dipoles directing toward the polymer in shell layer, the  $E_c$  of ZnO nanorod core is shifted toward the local vacuum level of the polymer by the interfacial dipole field,<sup>61</sup> along with up-shifting the vacuum level of ZnO nanorod core, while the position of quasi-Fermi level of electrons relative to the  $E_c$  of ZnO core is not affected (Figure 7b), as is similar to the interfacial modification with dipole organics in P3HT/TiO<sub>2</sub> and P3HT/ZnO solar cells.<sup>26-28</sup> The  $E_c$  of ZnO core is changed by  $e\delta E_c^{26,61}$  where  $\delta E$  is the change of surface potential and can be calculated from Poisson's equation,  $\delta E = N\mu\cos\theta/\varepsilon_r\varepsilon_0$ , where N is the dipole concentration,  $\mu$  the dipole moment,  $\theta$  the angle between the dipole and the surface normal,  $\varepsilon_{\mathrm{r}}$  the dielectric constant of ZnO, and  $\varepsilon_0$  the permittivity of free space. Therefore, the magnitude of  $E_{\rm c}$  shifting and thereby the device  $V_{\rm oc}$  correlates with the population of dipoles. More dipoles formed by thicker ZnO-QDs shell make a larger energy difference between the quasi-Fermi levels of electrons in ZnO nanorod core and holes in MEH-PPV for a higher  $V_{oc}$  in the range of L = 0-8 nm; however, increasing L from 8 up to 16 nm does not cause the remarkable  $E_c$  shift in ZnO nanorod core due to the saturated dipole formation and results in the almost unchanged  $V_{\rm oc}$ .

The interfacial dipole generation actually can be regarded as formation of a charger-transfer state of weakly bound pairs of electrons and holes within e—h separation of few nanometers by Coulombic attraction, <sup>59,64–67</sup> where the charges may either become free charge carriers after escaping their mutual attraction or recombine to loss energy. The charge recombination may reduce  $V_{\rm oc}$ .<sup>68–70</sup> Previous observations on HPSCs<sup>8,27,63,71–77</sup> or dye-sensitized solar cells (DSCs)<sup>78,79</sup> showed that the device with a larger  $V_{\rm oc}$  often exhibited a

reduced charge recombination rate (or longer  $\tau_e$ ), which seems to inform that  $V_{oc}$  is governed by  $\tau_e$  in both HPSCs<sup>11,16,18,62</sup> and DSCs.<sup>78,79</sup> Recent results have shown that there is no direct correlation between  $V_{oc}$  and  $\tau_e$  in HPSCs,<sup>13,17,34,47</sup> even though they may change in a similar trend subjected to the same influencing factor.<sup>13,34</sup> The *L*-dependence of  $\tau_e$  opposite to that of  $V_{oc}$  or  $V_{bi}$  further confirms no direct relationship between  $V_{oc}$  on  $\tau_e$  in MEH-PPV/ZnO-NR-QD-NA devices, for which the different effects on  $V_{oc}$  and  $\tau_e$  due to shell formation should be responsible, that is,  $V_{oc}$  is dominantly affected by the  $E_c$  shift ( $e\delta E$ ) of ZnO core due to dipole field, whereas  $\tau_e$  is mainly related to the interfacial charge recombination resulting from defect concentration.

3.2.4. Further Evidence from  $TiO_2$  Nanoarrays. To test whether the homogeneous core/shell strategy for improving  $V_{\rm oc}$  can be extended to other oxides, solar cells based on  $TiO_2$ -NA and  $TiO_2$ -NR-QD-NA were examined by the routine depicted in Figure 1. Insets to Figure 9 show the SEM images



**Figure 9.** J-V curves of TiO<sub>2</sub>-NA ( $\Box$ ) and TiO<sub>2</sub>-NR-QD-NA ( $\bigcirc$ ) solar cells under AM 1.5 illumination. Insets are the SEM images of TiO<sub>2</sub>-NA (I) and TiO<sub>2</sub>-NR-QD-NA (II).

of TiO2-NA and TiO2-NR-QD-NA. Clearly, the TiO2-NA consists of vertically aligned TiO<sub>2</sub> nanorods with a length of 500 nm and diameter of 40-50 nm (inset I to Figure 9); after TiO2-QDs deposition, the morphology of TiO2-NA is not remarkably altered, but the TiO<sub>2</sub> nanorod surfaces are decorated by nanoparticles (inset II to Figure 9). The XRD patterns of the TiO2-NA samples before and after TiO2-QDs deposition matched well rutile TiO<sub>2</sub> (see Figure S6 in the Supporting Information). TEM images showed that the original TiO<sub>2</sub> nanorods had quite smooth surfaces, and a core/shell structure was formed after deposition TiO<sub>2</sub>-QDs (see Figure S7 in the Supporting Information). HRTEM images (see Figure S7 in the Supporting Information) indicated that the TiO<sub>2</sub> nanorod grew along [001] direction, and the shell consisted of differently oriented TiO<sub>2</sub>-QDs of 2-4 nm in size. Therefore, the TiO<sub>2</sub>-NR-QD-NA with the single-crystalline TiO<sub>2</sub> nanorod as core and the polycrystalline film (5 nm in thickness) of TiO<sub>2</sub>-QDs as shell is successfully by the combined hydrothermal-solvothermal method. Similar to the cases of ZnO nanoarrays, MEH-PPV can also be well infiltrated into the interspaces between either TiO2-NA or TiO2-NR-QD-NA nanorods after thermal annealing (150 °C, 20 min) (see Figure S8 in the Supporting Information). The J-V curves of MEH-PPV/TiO2-NA and MEH-PPV/TiO2-NR-QD-NA solar cells under AM1.5 illumination are compared in Figure 9. The MEH-PPV/TiO<sub>2</sub>-NA device exhibits the  $V_{oc}$  of 0.35 V,  $J_{sc}$  of 3.54 mA/cm<sup>2</sup>, FF of 37.70% and  $\eta$  of 0.47%, whereas the MEH-PPV/TiO<sub>2</sub>-NR-QD-NA device shows the  $V_{oc}$  of 0.71 V,  $J_{sc}$  of 4.85 mA/cm<sup>2</sup>, FF of 46.62% and  $\eta$  of 1.60%. The observed  $V_{\rm oc}$ of 0.35 V in MEH-PPV/TiO2-NA device is comparable to reports on P3HT/TiO2-NA solar cells.<sup>11</sup> The phenomena similar to the cases of ZnO are observed upon the TiO2-QDs shell formation, including the significant increase in  $V_{oc}$  as well as  $J_{sc}$ , the greatly reduced  $\mu_e$  but much higher  $V_{bi}$  ( $\mu_e = 9.6 \times 10^{-9} \text{ m}^2 \text{v}^{-1} \text{s}^{-1}$  and  $V_{bi} = 0.42 \text{ V}$  for MEH-PPV/TiO<sub>2</sub>-NA,  $\mu_e = 3.6 \times 10^{-11} \text{ m}^2 \text{ V}^{-1} \text{ s}^{-1}$  and  $V_{bi} = 0.73 \text{ V}$  and for MEH-PPV/ TiO<sub>2</sub>-NR-QD-NA composites) (see Figure S9 in the Supporting Information). IMVS results (see Figure S10 in the Supporting Information) showed also that the significant reduction in  $\tau_e$  upon the shell formation from 6.78 to 1.50 ms, confirming no direct relationship between  $V_{\rm oc}$  on  $\tau_{\rm e}$  in MEH-PPV/TiO<sub>2</sub>-NR-QD-NA devices. Note, the observed  $\tau_{a}$  in our devices is comparable to the reciprocal of the rate constants for recombination between the electrons in TiO<sub>2</sub> and holes in P3HT determined by transient photovoltage method.<sup>27,77,80</sup> The much larger  $\mu_e$  in MEH-PPV/ZnO-NA (1.8 × 10<sup>-6</sup> m<sup>2</sup>  $V^{-1} s^{-1}$ ) than in MEH-PPV/TiO<sub>2</sub>-NA (9.6 × 10<sup>-9</sup> m<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>) agrees with the fact that electron mobility in ZnO is a least 1-2order of magnitude higher than in TiO<sub>2</sub>.<sup>4,80</sup> Moreover, the FF increases by formation of TiO2 homogeneous core/shell structure, which is similar to the case of ZnO (Table 1).

#### 4. CONCLUSION

We develop a strategy to significantly increase the  $V_{\rm oc}$  in polymer/oxide-nanoarray solar cells by formation of homogeneous core/shell nanoarray, in which single-crystalline nanorod serving as core and the film of corresponding quantum dots acts as shell. The  $V_{\rm oc}$  in the hybrid solar cells based on the core/ shell nanoarrays is dominantly governed by the quasi-Fermi levels of the electrons in the core and holes in the polymer. The shell formation reduces the mobility difference between the photogenerated electrons in the core and holes in the polymer, creating a dipole layer between the core and the polymer. Upon the dipole layer formation, the  $E_c$  edge of core is shifted toward the local vacuum level of the polymer with an increase in the energy difference between the quasi-Fermi levels of the electrons and holes for a higher  $V_{\rm oc}$  until the saturated dipole formation at the shell thickness that is comparable to the threshold e-h separation for electronic coupling. Moreover, our results indicate that increasing  $V_{\rm oc}$  by the core/shell strategy seems to be not dependent on the kinds of metal oxides.

## ASSOCIATED CONTENT

## Supporting Information

XRD patterns of ZnO-NA and ZnO-NR-QD-NAs; SEM and TEM (HRTEM) images of ZnO-NR-QD-NAs with a time of t = 4 and 6 h; fitting experimental data of ZnO-NA and ZnO-NR-QD-NAs based devices by SCLC method; calculated dependence of ZnO conduction band edge shift  $\Delta E_C$  on quantum dot size; XRD patterns of TiO<sub>2</sub>-NA and TiO<sub>2</sub>-NR-QD-NA; TEM and HRTEM images of TiO<sub>2</sub>-NA and TiO<sub>2</sub>-NR-QD-NA; fitting experimental data of TiO<sub>2</sub>-NA and TiO<sub>2</sub>-NR-QD-NA based devices by SCLC method; IMVS spectra of solar cells based on TiO<sub>2</sub>-NA and TiO<sub>2</sub>-NR-QD-NA. This material is available free of charge via the Internet at http:// pubs.acs.org.

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F.W. performed the experiments, and drafted the manuscript. Q.C., Z.Q., C.L., and H.Z. participated in the theoretical analysis. W.S. participated in the device preparation. M.W. conceived of the study, and participated in its design and coordination. All authors have read and approved the final manuscript.

#### Notes

The authors declare no competing financial interest.

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